

Introduction to Condensation Heat Transfer

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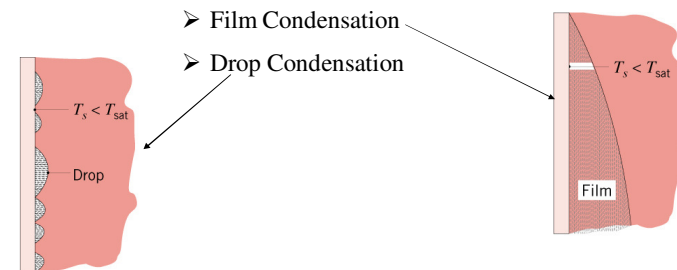
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Condensation Heat Transfer-I

- Condensation implies transformation of vapor back to liquid
- There are basically two mechanisms for condensation



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Condensation Heat Transfer-II

- Film Condensation
 - Entire surface is covered by the **condensate**, which flows continuously from the surface and provides a resistance to heat transfer between the vapor and the surface.
 - Thermal resistance is reduced through use of short vertical surfaces and horizontal cylinders.
- Dropwise Condensation
 - Surface is covered by drops ranging from a few micrometers to agglomerations visible to the naked eye.
 - Thermal resistance is greatly reduced due to absence of a continuous film.
 - Surface coatings may be applied to inhibit *wetting* and stimulate dropwise condensation.

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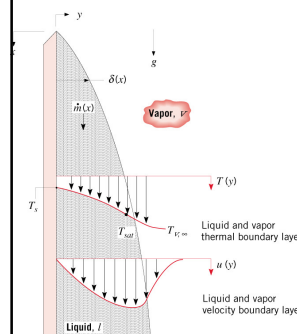
Analysis of Film Condensation-I

Nusselt Analysis for Laminar Flow

Assumptions:

- A pure vapor at T_{sat}
- Creeping boundary layer flow
- Negligible shear stress at liquid/vapor interface.

We shall derive it from the basic equations



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Analysis of Film Condensation-II

- Momentum Equation

$$\rho_l \left(u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} \right) = - \frac{\partial p}{\partial x} + \mu_l \left(\frac{\partial^2 u}{\partial y^2} \right) + \rho_l g$$

Negligible

 $\rho_v g$

$$\Rightarrow (\rho_l - \rho_v)g = -\mu_l \left(\frac{d^2 u}{dy^2} \right) \Rightarrow \left(\frac{d^2 u}{dy^2} \right) = - \frac{(\rho_l - \rho_v)g}{\mu_l}$$

On Integration we get

$$u = - \frac{(\rho_l - \rho_v)g}{\mu_l} \frac{y^2}{2} + c_1 y + c_2$$

Boundary Conditions

$$u = 0, \quad y = 0 \Rightarrow c_2 = 0$$

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Analysis of Film Condensation-III

$$\frac{\partial u}{\partial y} = 0, \quad y = \delta$$

$$\Rightarrow c_1 = \frac{\rho_l - \rho_v}{\mu_l} g \delta$$

$$\Rightarrow u = \frac{(\rho_l - \rho_v)g\delta^2}{\mu_l} \left(\frac{y}{\delta} - \frac{y^2}{2\delta^2} \right)$$

- Energy Equation

$$\left(u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} \right) = \alpha \left(\frac{\partial^2 T}{\partial y^2} \right) \Rightarrow \left(\frac{\partial^2 T}{\partial y^2} \right) = 0$$

Negligible (Advection neglected)

$$\Rightarrow T = c_1 y + c_2$$

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Analysis of Film Condensation-IV

Boundary Conditions

$$T = T_w, \quad y = 0$$

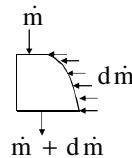
$$T = T_{sat}, \quad y = \delta$$

$$\Rightarrow T = T_w + (T_{sat} - T_w) \frac{y}{\delta}$$

Note that the boundary layer thickness is still unknown

This is obtained in the following manner

$$\begin{aligned} \dot{m} &= \int_0^\delta \rho_l u dy \quad \text{For unit width} \\ &= \int_0^\delta \rho_l \frac{(\rho_l - \rho_v)g\delta^2}{\mu_l} \left(\frac{y}{\delta} - \frac{y^2}{2\delta^2} \right) dy \\ &= \rho_l \frac{(\rho_l - \rho_v)g\delta^3}{3\mu_l} \end{aligned}$$



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Analysis of Film Condensation-V

$$d\dot{m} = \rho_l \frac{(\rho_l - \rho_v)g3\delta^2}{3\mu_l} d\delta \quad (1)$$

From Energy Balance per unit width

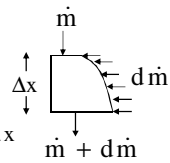
Advection neglected implies that heat travels only in y direction

$$h_{fg} d\dot{m} = \left(k_l \frac{\partial T}{\partial y} \Big|_{y=\delta} \right) \Delta x = k_l \left(\frac{T_{sat} - T_w}{\delta} \right) \Delta x$$

$$d\dot{m} = \frac{k_l}{h_{fg}} \left(\frac{T_{sat} - T_w}{\delta} \right) \Delta x \quad (2)$$

$$\text{Equating 1 and 2 we get} \quad \frac{\rho_l(\rho_l - \rho_v)g\delta^2}{\mu_l} d\delta = \frac{k_l}{h_{fg}} \left(\frac{T_{sat} - T_w}{\delta} \right) \Delta x$$

$$\Rightarrow \delta^3 \frac{d\delta}{dx} = \frac{k_l \mu_l (T_{sat} - T_w)}{\rho_l (\rho_l - \rho_v) h_{fg} g} \Rightarrow \frac{\delta^4}{4} = \frac{k_l \mu_l (T_{sat} - T_w)}{\rho_l (\rho_l - \rho_v) h_{fg} g} x + c_1$$



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Analysis of Film Condensation-VI

Using the Boundary Condition $\delta = 0$ at $x = 0 \Rightarrow c_1 = 0$

$$\Rightarrow \delta = \left(\frac{4k_1\mu_1(T_{\text{sat}} - T_w)x}{\rho_1(\rho_1 - \rho_v)h_{fg}g} \right)^{1/4}$$

To account for sub-cooling the following correction is made

$$h'_{fg} = h_{fg}(1 + 0.68 \text{Ja})$$

$$\text{Where, } \text{Ja} = \frac{c_p(T_{\text{sat}} - T_w)}{h_{fg}} \quad \text{Jacob Number}$$

Finally,

$$h = \frac{q''}{T_{\text{sat}} - T_w} = \frac{1}{T_{\text{sat}} - T_w} k_1 \frac{\partial T}{\partial y} = \frac{k_1}{T_{\text{sat}} - T_w} \frac{(T_{\text{sat}} - T_w)}{\delta} = \frac{k_1}{\delta}$$

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Analysis of Film Condensation-VII

$$\Rightarrow \text{Nu}_x = \frac{hx}{k_1} = \frac{k_1}{\delta} \frac{x}{k_1} = \frac{x}{\delta} \Rightarrow \text{Nu}_x = \left(\frac{\rho_1(\rho_1 - \rho_v)h'_{fg}gx^3}{k_1\mu_1(T_{\text{sat}} - T_w)} \right)^{1/4}$$

$$\text{Nu}_x = 0.707 \left(\frac{\rho_1(\rho_1 - \rho_v)h'_{fg}gx^3}{k_1\mu_1(T_{\text{sat}} - T_w)} \right)^{1/4}$$

By integration one can show that

$$\overline{\text{Nu}}_L = \frac{4}{3} \text{Nu}_L = 0.943 \left(\frac{\rho_1(\rho_1 - \rho_v)h'_{fg}gL^3}{k_1\mu_1(T_{\text{sat}} - T_w)} \right)^{1/4}$$

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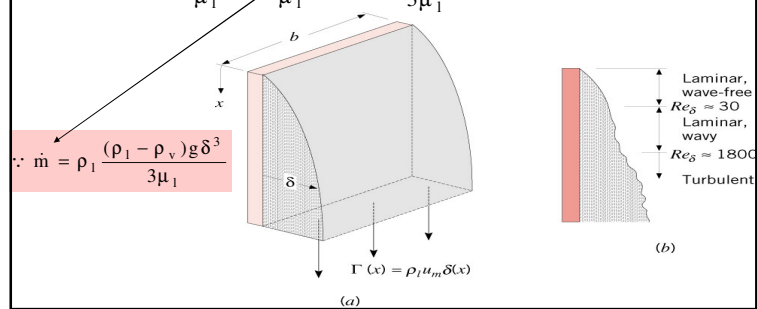
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Analysis of Film Condensation-VIII

• Effects of Turbulence

- Three flow regimes may be identified and delineated in terms of a Reynolds number defined as

$$\Rightarrow \text{Re}_\delta = \frac{4\rho_1\bar{u}\delta}{\mu_1} = \frac{4\dot{m}}{\mu_1} = \frac{4g\rho_1(\rho_1 - \rho_v)\delta^3}{3\mu_1^2}$$



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Analysis of Film Condensation-IX

Laminar Region $\text{Re}_\delta < 30$

$$\frac{\bar{h}_L(v_1^2/g)^{1/3}}{k_1} = 1.47 (\text{Re}_\delta)^{-1/3}$$

Wavy Laminar Region $30 < \text{Re}_\delta < 1800$

$$\frac{\bar{h}_L(v_1^2/g)^{1/3}}{k_1} = \frac{\text{Re}_\delta}{1.08 (\text{Re}_\delta)^{1.22} - 5.2}$$

Turbulent Region $\text{Re}_\delta > 1800$

$$\frac{\bar{h}_L(v_1^2/g)^{1/3}}{k_1} = \frac{\text{Re}_\delta}{8750 + 58 \text{Pr}_1^{-0.5} (\text{Re}_\delta^{0.75} - 253)}$$

We can obtain $\bar{h}_L$ provided we know Re_δ

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Analysis of Film Condensation-X

➤ Calculation procedure:

To use the equations in the previous slide, we first use the relation for mass flow rate per unit width

$$\dot{m}_{(x=L)} = \frac{\bar{h}_L \times L (T_{\text{sat}} - T_w)}{h'_{fg}}$$

$$\text{Since } \text{Re}_\delta = \frac{4 \dot{m}}{\mu_l} = \frac{4}{\mu_l} \frac{\bar{h}_L L (T_{\text{sat}} - T_w)}{h'_{fg}}$$

$$\Rightarrow \bar{h}_L = \frac{\text{Re}_\delta h'_{fg} \mu_l}{4 L (T_{\text{sat}} - T_w)}$$

Thus, the left hand side of the equations in previous slide can also be expressed in terms of Reynolds number and known parameters. Hence we can solve for Re_δ , and hence for $\bar{h}_L$

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Analysis of Film Condensation-XI

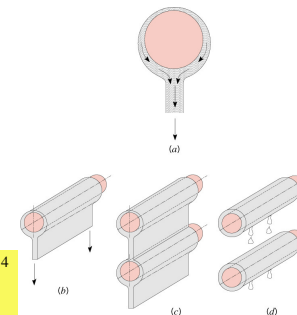
- A single tube or sphere:

$$\bar{h}_D = C \left(\frac{\rho_l (\rho_l - \rho_v) k_l^3 h'_{fg} g}{\mu_l (T_{\text{sat}} - T_w) D} \right)^{1/4}$$

Tube: $C = 0.729$ Sphere: $C = 0.826$

- Vertical Tier of N-tubes:

$$\bar{h}_{D,N} = 0.729 \left(\frac{\rho_l (\rho_l - \rho_v) k_l^3 h'_{fg} g}{N \mu_l (T_{\text{sat}} - T_w) D} \right)^{1/4}$$



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Condensation Inside Tubes

- Shah Correlation (1979):

$$h_f(x) \frac{D}{k_f} = 0.023 \text{Re}_e^{0.8} \text{Pr}_f^{0.4} f(x, p_R) \quad p_R = p/p_c$$

$$f(x, p_R) = \left[(1-x)^{0.8} + (3.8x^{0.76} (1-x)^{0.04}) p_R^{-0.38} \right]$$

It is suggested that $x = 0.999$ to start the method properly rather than $x = 1$.

Applicability in the range

$$10.8 < G < (\text{kg/m}^2\text{-s})$$

$$\text{Re}_{f0} > 350 \text{ for tubes and } \text{Re}_{f0} > 3000 \text{ for annuli}$$

$$0.002 < p_R < 0.44$$

$$0 < x < 1$$