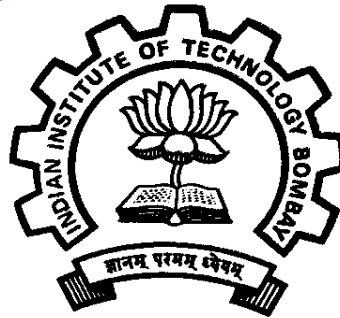


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ME 209

Basic Thermodynamics (Lecture-12)

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Review of Lecture 11

- Analysed a Carnot engine and showed that the Thermodynamic scale and Ideal Gas Kelvin scale are identical
- Derived the Clausius Inequality, defined entropy and entropy production
- Derived the Second law in terms of change of entropy for an infinitesimal process
- Stated the principle of increase in entropy of the universe.

Agenda for Today

- Understand property relations
- Understand h-s and T-s diagrams
- Derive the Second law for open systems

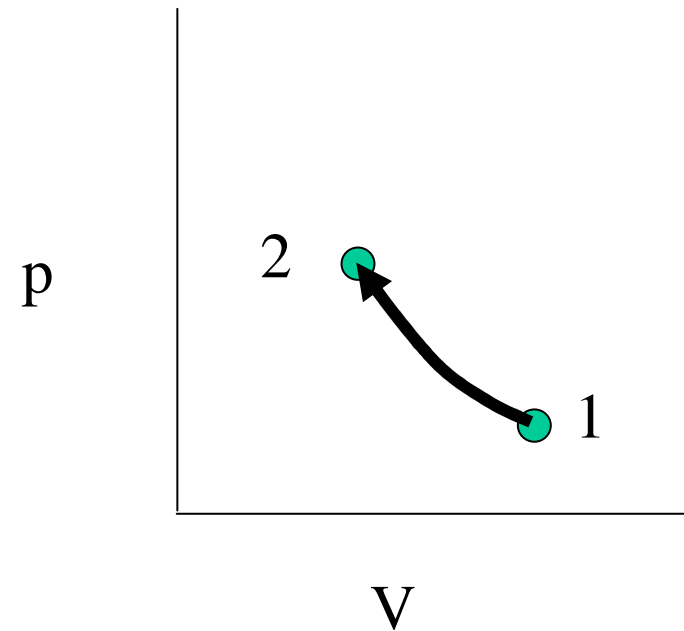
TdS Relations-I

- Consider two states 1 and 2 as shown
- The entropy change $S_2 - S_1$ shall not depend on path
- Let us identify the reversible path 1-2
- First law for the closed system can be written as

$$dQ = dU + dW$$

$$\Rightarrow TdS = dU + pdV$$

Note that the above relation involves only properties and hence shall be independent of path



TdS Relations-II

- We can also modify the relation derived as follows

$$TdS = dU + pdV$$

$$\Rightarrow TdS = dU + d(pV) - Vdp$$

$$\Rightarrow TdS = dH - Vdp$$

- The above two property relations are the working equations for the evaluation of entropy change and are known as TdS relations

T-S Diagram-I

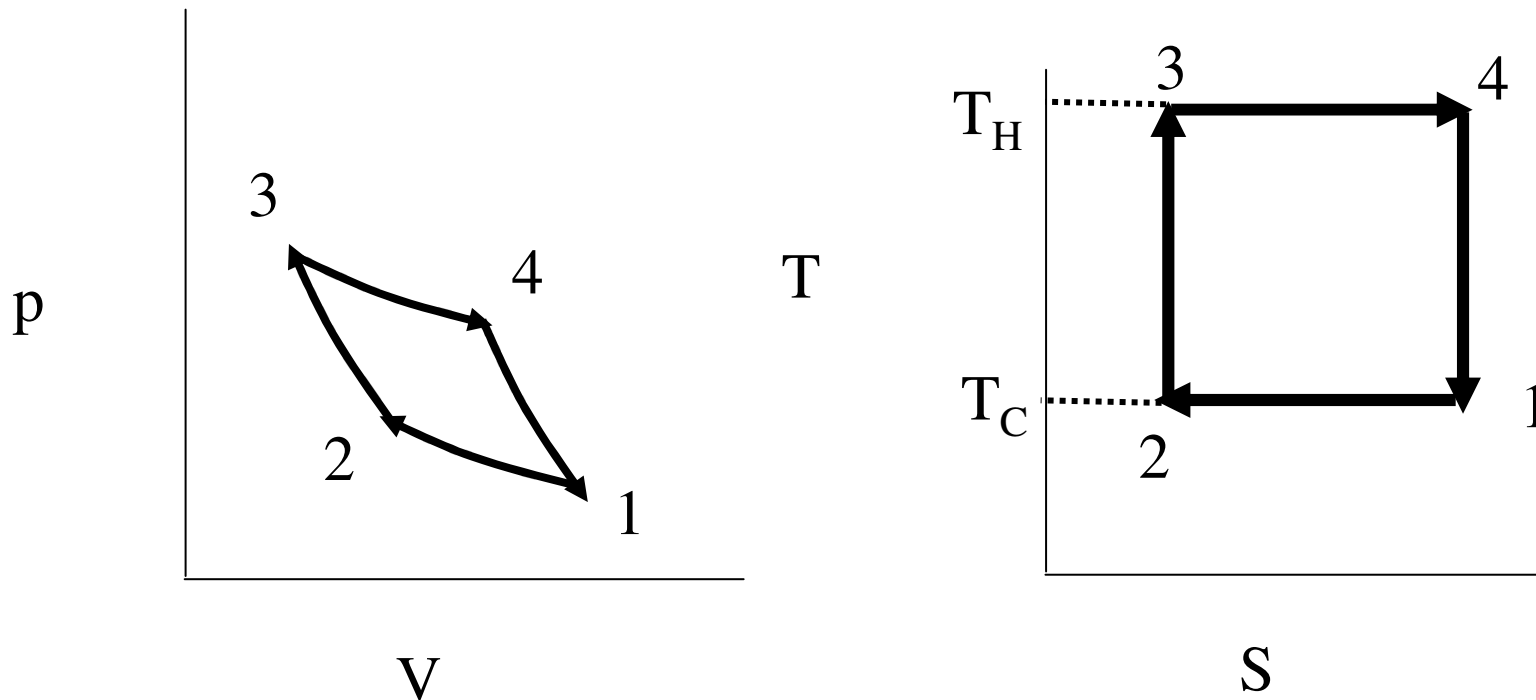
- For adiabatic and reversible process $ds = 0$
- Further, for a reversible process, the heat transferred per unit mass of the substance can be written as

$$\Rightarrow Q_R = \int_1^2 T ds$$

- The above has motivated construction of T-s diagram. Its usefulness is established by considering analysis of Carnot Cycle

T-S Diagram-II

1. Isothermal Compression (1-2)
2. Adiabatic Compression (2-3)
3. Isothermal Expansion (3-4)
4. Adiabatic Expansion (4-1)



T-S Diagram-III

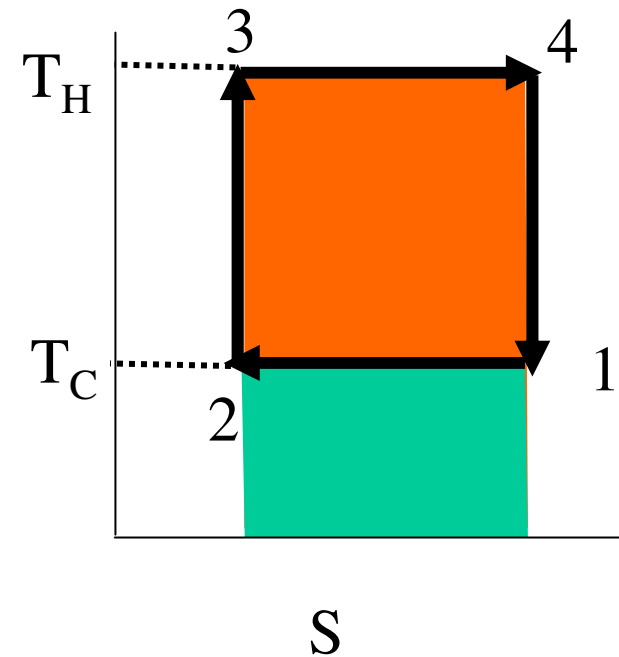
- Heat rejected in process 1-2

$$\Rightarrow Q_C = -Q_{12} = -\int_1^2 T ds = -T_C (S_2 - S_1)$$

- Heat added in process 3-4

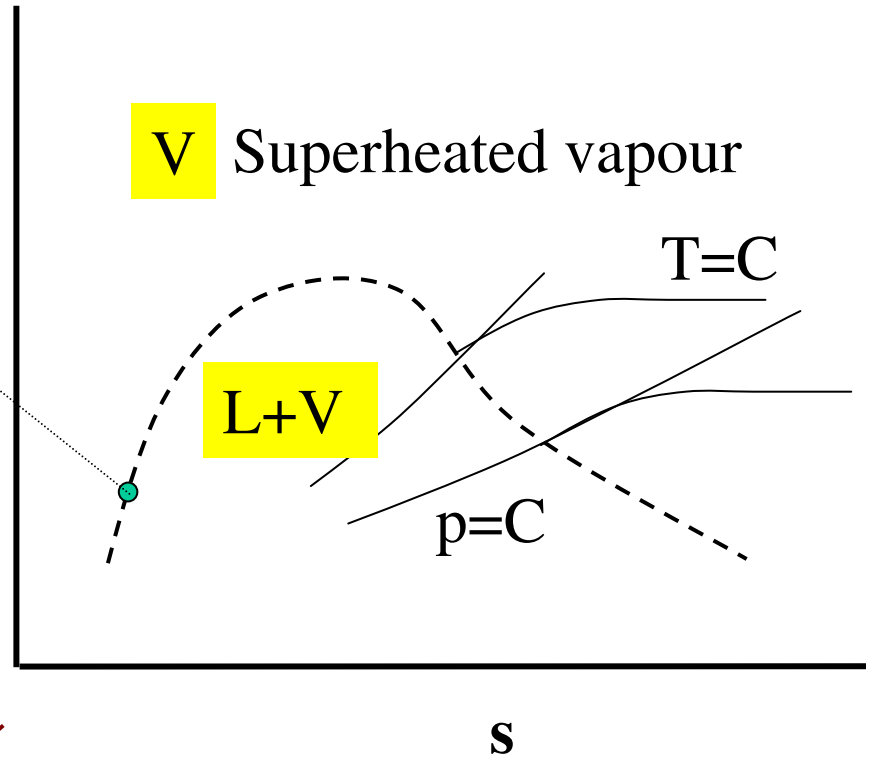
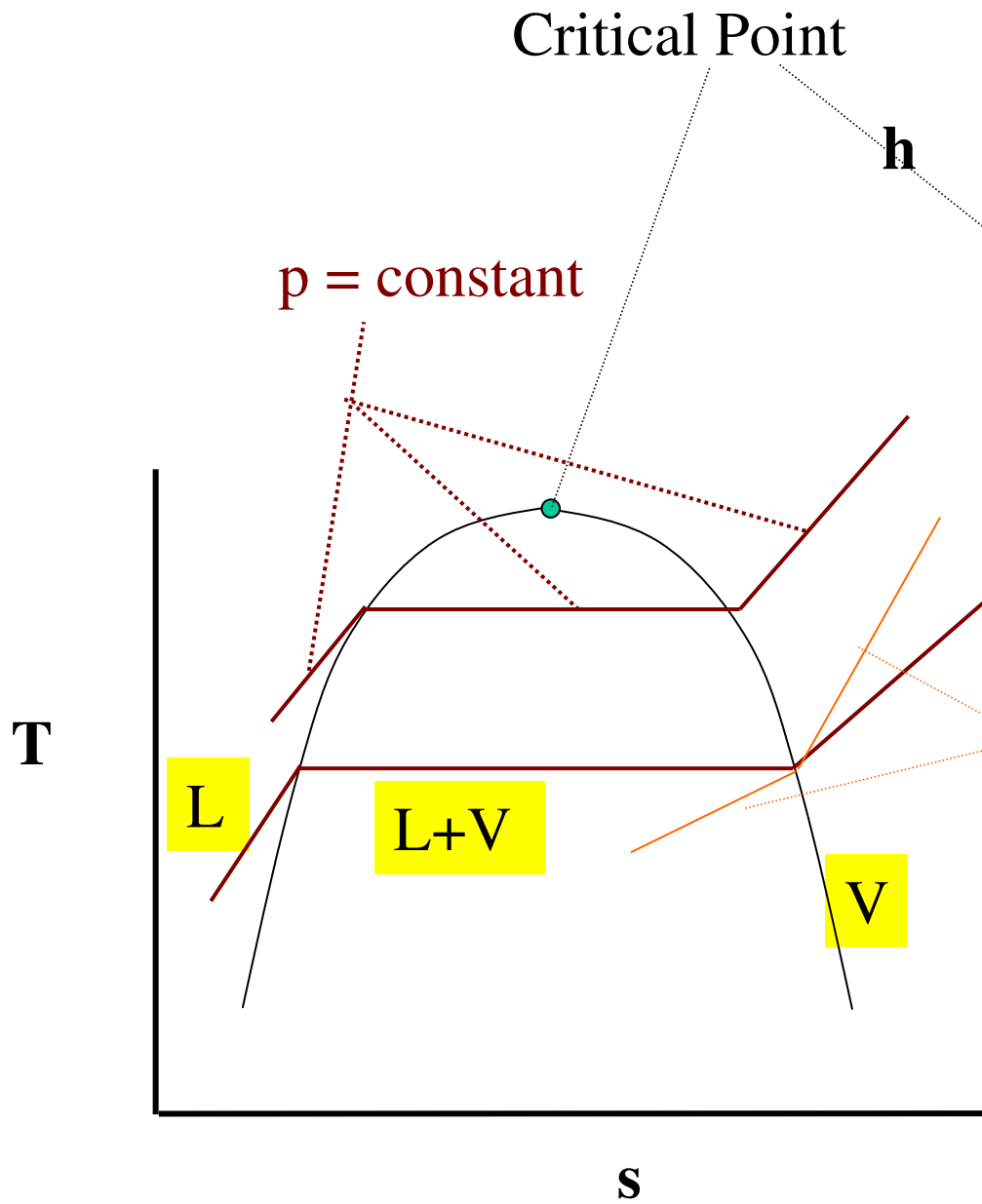
$$\Rightarrow Q_H = Q_{34} = \int_3^4 T ds = T_H (S_4 - S_3)$$

$$\begin{aligned} \Rightarrow \eta &= 1 - \frac{Q_C}{Q_H} \\ &= 1 - \frac{T_C (S_1 - S_2)}{T_H (S_4 - S_3)} = 1 - \frac{T_C}{T_H} \end{aligned}$$



h-s diagram

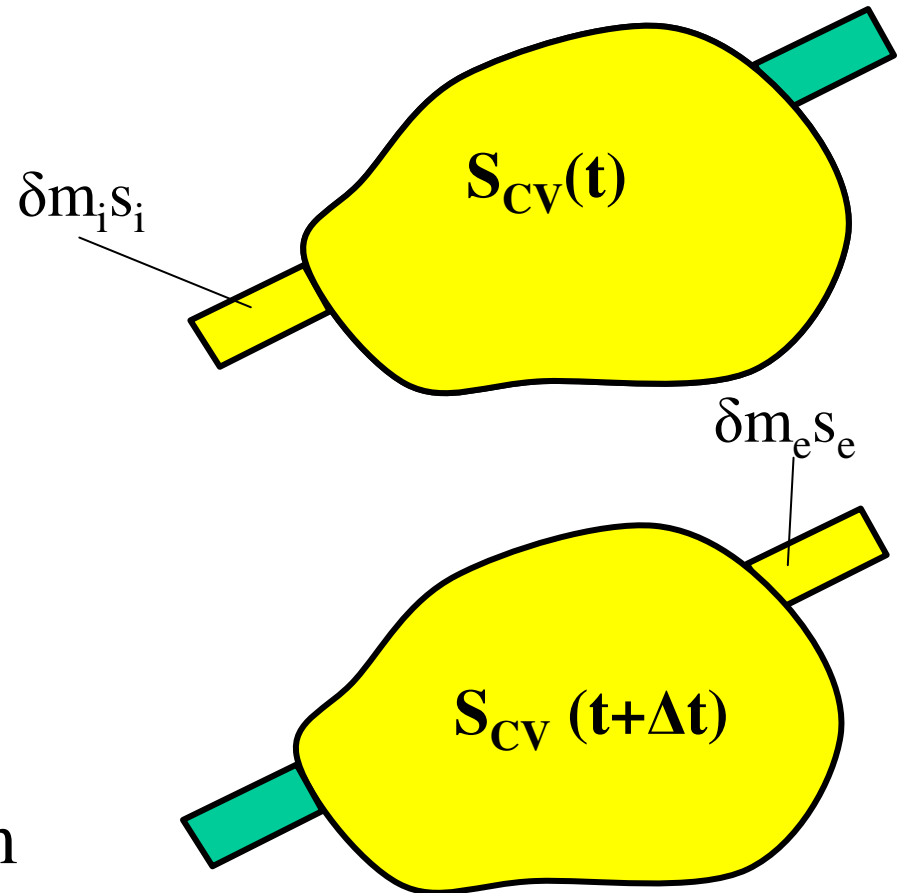
- For open systems, with kinetic and potential effects neglected, we saw that for turbines, pumps/compressor and heat exchangers the heat/work per unit mass of the substance was directly related to change of enthalpy
- This motivated inventions of h-s diagrams, particularly for water substance
- These are called Mollier Diagrams



$v = \text{constant}$

II Law for Open System-I

- Now let us consider the arbitrary system again
- At time, t , the total entropy of the control mass is $S_{CV}(t) + \delta m_i s_i$
- At time, $t + \Delta t$, the total entropy of the control mass is $S_{CV}(t + \Delta t) + \delta m_e s_e$
- During this Δt , the system has heat interaction of δQ . Usual Sign convention applies



II Law for Open System-II

- The second law for the control mass can be written as,

$$S_{CV}(t+\Delta t) + \delta m_e s_e - (S_{CV}(t) + \delta m_i s_i) = (\delta Q/T) + \delta S_p$$

- Rearranging and dividing both sides by Δt and shrinking the Δt to 0, we can write,

$$\left. \frac{S_{CV}(t+\Delta t) - S_{CV}(t)}{\Delta t} \right|_{\Delta t \rightarrow 0} = \left. \frac{\delta Q}{\Delta t} \frac{1}{T} \right|_{\Delta t \rightarrow 0} + s_i \left. \frac{\delta m_i}{\Delta t} \right|_{\Delta t \rightarrow 0} - s_e \left. \frac{\delta m_e}{\Delta t} \right|_{\Delta t \rightarrow 0} + \left. \frac{\delta S_p}{\Delta t} \right|_{\Delta t \rightarrow 0}$$

$$\text{or, } \frac{dS_{CV}}{dt} = \frac{\dot{Q}_{CV}}{T} + \dot{m}_i s_i - \dot{m}_e s_e + \dot{S}_p$$

Evaluation of Change in Entropy-I

For the system with ideal gas

- Once the two states are known we can use the TdS relations to evaluate the entropy change

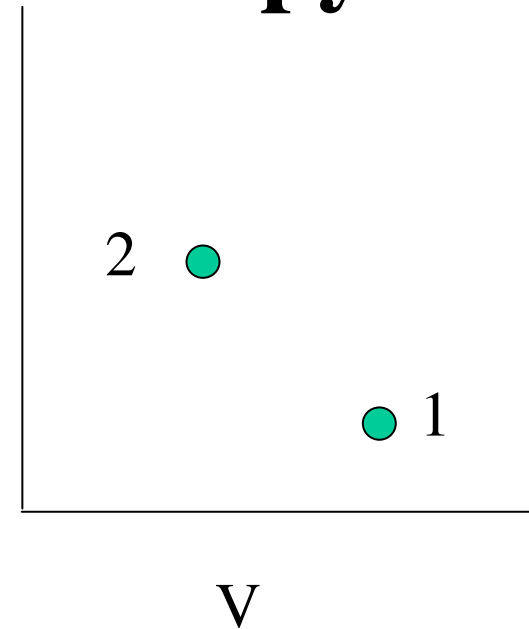
$$TdS = dU + pdV$$

$$\Rightarrow TdS = mc_v dT + \frac{mRT}{V} dV$$

$$\Rightarrow dS = mc_v \frac{dT}{T} + mR \frac{dV}{V}$$

$$\Rightarrow \int_1^2 dS = mc_v \int_1^2 \frac{dT}{T} + mR \int_1^2 \frac{dV}{V}$$

$$\Rightarrow S_2 - S_1 = mc_v \ln \frac{T_2}{T_1} + mR \ln \frac{V_2}{V_1}$$



Evaluation of Change in Entropy-II

- Proceeding on similar lines, we can use the second TdS relation

$$TdS = dH - Vdp \quad \text{and show}$$

$$\Rightarrow S_2 - S_1 = mc_p \ln \frac{T_2}{T_1} - mR \ln \frac{p_2}{p_1}$$

Evaluation of Change in Entropy-III

For the system with incompressible substance

- For incompressible substances the expansion work is absent that makes $c_p = c_v = c$

$$TdS = dU + pdV$$

$$\Rightarrow TdS = dU + pdV$$

$$\Rightarrow TdS = mcdT$$

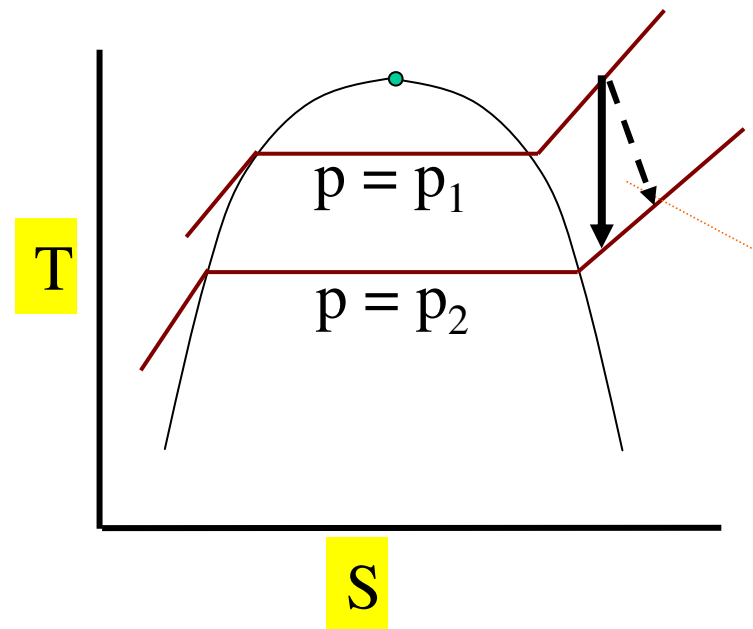
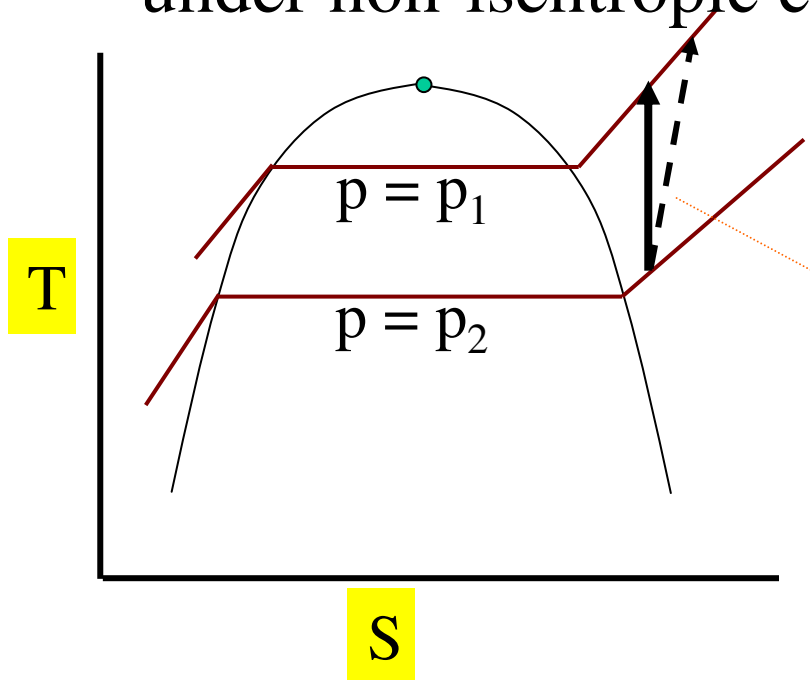
$$\Rightarrow dS = mc \frac{dT}{T}$$

$$\Rightarrow \int_1^2 dS = mc \int_1^2 \frac{dT}{T}$$

$$\Rightarrow S_2 - S_1 = mc \ln \frac{T_2}{T_1}$$

Isentropic efficiency-I

- Turbines, compressors and nozzles, perform ideally under isentropic conditions
- Turbine output decreases, Compressor input increases and nozzle KE developed decreases, under non-isentropic conditions



Isentropic efficiency-II

- We define isentropic efficiency as

$$\eta_{s-\text{Turbine/Nozzle}} = \frac{h_i - h_e}{h_i - h_{e^*}}$$

$$\eta_{s-\text{comp/pump}} = \frac{h_{e^*} - h_i}{h_e - h_i}$$