

^{3/15} Concept of Exact Differential-II
$dv_{1-2} = \frac{\partial v}{\partial p} \bigg _{T1} dp + \frac{\partial v}{\partial T} \bigg _{p2} dT = \frac{\partial v}{\partial T} \bigg _{p1} dT + \frac{\partial v}{\partial p} \bigg _{T2} dp$
$\left(\frac{\partial \mathbf{v}}{\partial \mathbf{p}}\Big _{\mathrm{T1}} \mathrm{d}\mathbf{p} + \frac{\partial \mathbf{v}}{\partial \mathrm{T}}\Big _{\mathrm{p2}} \mathrm{d}\mathrm{T}\right) - \left(\frac{\partial \mathbf{v}}{\partial \mathrm{T}}\Big _{\mathrm{p1}} \mathrm{d}\mathrm{T} + \frac{\partial \mathbf{v}}{\partial \mathrm{p}}\Big _{\mathrm{T2}} \mathrm{d}\mathbf{p}\right) = 0$
$\Rightarrow \left(\frac{\partial \mathbf{v}}{\partial \mathbf{p}}\Big _{\mathbf{T}^2} - \frac{\partial \mathbf{v}}{\partial \mathbf{p}}\Big _{\mathbf{T}^1}\right) d\mathbf{p} = \left(\frac{\partial \mathbf{v}}{\partial \mathbf{T}}\Big _{\mathbf{p}^2} - \frac{\partial \mathbf{v}}{\partial \mathbf{T}}\Big _{\mathbf{p}^1}\right) d\mathbf{T}$
$\Rightarrow \frac{\left(\frac{\partial \mathbf{v}}{\partial \mathbf{p}}\Big _{\mathbf{T}_{2}} - \frac{\partial \mathbf{v}}{\partial \mathbf{p}}\Big _{\mathbf{T}_{1}}\right)}{\mathbf{dT}} = \frac{\left(\frac{\partial \mathbf{v}}{\partial \mathbf{T}}\Big _{\mathbf{p}^{2}} - \frac{\partial \mathbf{v}}{\partial \mathbf{T}}\Big _{\mathbf{p}^{1}}\right)}{\mathbf{dp}}$

Concept of Exact Differential-III

• As the points are close by, we can write the previous equation as

$$\Rightarrow \frac{\partial}{\partial T} \left(\frac{\partial v}{\partial p} \Big|_{T} \right)_{p} = \frac{\partial}{\partial p} \left(\frac{\partial v}{\partial T} \Big|_{p} \right)_{T}$$
$$\Rightarrow \frac{\partial^{2} v}{\partial T \partial p} = \frac{\partial^{2} v}{\partial p \partial T}$$

- The value of a mixed differential is independent of the order of differentiation
- The result is a consequence of assuming dv is independent of direction. Those differentials that satisfy this property are called exact differentials

^{5/15} Concept of Exact Differential-IV

- Every property change is an exact differential and Every exact differential represents change of a property
- The whole thing can be generalised as, given three variables x, y and z and they have a relation of the form

dz = M(x, y)dx + N(x, y)dy

then the differential dz is exact, if



• In thermodynamics, knowing that properties are exact, we shall equate the cross derivatives

Rules of Partial Derivatives-I

• Now let us look at the relation between partial derivatives

$$dx = \frac{\partial x}{\partial y} \left| dy + \frac{\partial x}{\partial z} \right|_{y} dz$$

• Similarly, we can write

6/15

8/15

$$dy = \frac{\partial y}{\partial x} \bigg|_{z} dx + \frac{\partial y}{\partial z} \bigg|_{x} dz$$

• Substituting the expression for dy in second equation into the first equation, we get

$$dx = \frac{\partial x}{\partial y} \bigg|_{z} \left(\frac{\partial y}{\partial x} \bigg|_{z} dx + \frac{\partial y}{\partial z} \bigg|_{x} dz \right) + \frac{\partial x}{\partial z} \bigg|_{y} dz$$

^{7/15} **Rules of Partial Derivatives-II** • Collecting the coefficients of dx and dz, we can write $dx \left(1 - \frac{\partial x}{\partial y} \Big|_{z} \frac{\partial y}{\partial x} \Big|_{z} \right) = \left(\frac{\partial x}{\partial y} \Big|_{z} \frac{\partial y}{\partial z} \Big|_{x} + \frac{\partial x}{\partial z} \Big|_{y} \right) dz$ • Now if we go to two neighboring states such that dz = 0 and dx \neq 0, then it is necessary to have $1 - \frac{\partial x}{\partial y} \Big|_{z} \frac{\partial y}{\partial x} \Big|_{z} = 0 \qquad \text{Or } \frac{\partial x}{\partial y} \Big|_{z} \frac{\partial y}{\partial x} \Big|_{z} = 1 \qquad \text{Or } \frac{\partial x}{\partial y} \Big|_{z} = \frac{1}{\frac{\partial y}{\partial x} \Big|_{z}}$ • We can call the above as Reciprocal Rule

Rules of Partial Derivatives-III

• Now if we go to two neighboring states such that dx = 0and $dz \neq 0$, then, we can write,

$$\frac{\partial x}{\partial y} \bigg|_{z} \frac{\partial y}{\partial z} \bigg|_{x} = -\frac{\partial x}{\partial z} \bigg|_{y}$$

• Now applying the reciprocal rule, we can write

$$\frac{\partial \mathbf{x}}{\partial \mathbf{y}}\Big|_{\mathbf{z}} \frac{\partial \mathbf{y}}{\partial \mathbf{z}}\Big|_{\mathbf{x}} = \frac{-1}{\frac{\partial \mathbf{z}}{\partial \mathbf{x}}\Big|_{\mathbf{y}}} \quad \text{Or } \frac{\partial \mathbf{x}}{\partial \mathbf{y}}\Big|_{\mathbf{z}} \frac{\partial \mathbf{y}}{\partial \mathbf{z}}\Big|_{\mathbf{x}} \frac{\partial \mathbf{z}}{\partial \mathbf{x}}\Big|_{\mathbf{y}} = -1$$

• The above can be called as Cyclic Rule

Thermodynamic Functions-I

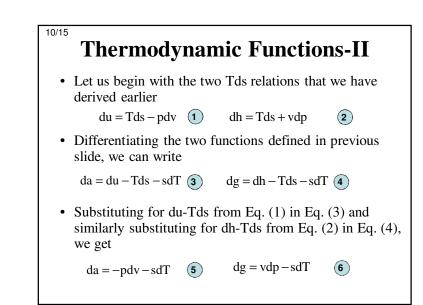
- Having laid the foundation for deriving property relations, we shall now go on to define two new thermodynamic functions called Helmholtz and Gibbs functions
- The Helmholtz function A and Gibbs function G are defined as
 - A = U TS G = H TS

9/15

• Their intensive counterparts are **a** and **g**

a = u - Ts g = h - Ts

• Now we shall begin manipulating these functions



^{11/15} **Thermodynamic Functions-III** • Now let us manipulate Eq. (1), where u = u(s,v) du = Tds - pdv• Chain rule implies $du = \frac{\partial u}{\partial s}\Big|_{v} ds + \frac{\partial u}{\partial v}\Big|_{s} dv$ • Comparing the above two equations, we can write, $\frac{\partial u}{\partial s}\Big|_{v} = T, \quad \frac{\partial u}{\partial v}\Big|_{s} = -p$ • From Eq. (2) dh = Tds + vdp• We can write, $\frac{\partial h}{\partial s}\Big|_{s} = T, \quad \frac{\partial h}{\partial p}\Big|_{s} = v$ (8)

